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Bathymetry enhancement by altimetry-derived gravity anomalies in the East Sea (Sea of Japan)

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Abstract The gravity-geologic method (GGM) was used to enhance the bathymetry of the East Sea (Sea of Japan) with satellite altimetry-derived free-air gravity anomalies and shipborne depth measurements. By comparison with the bathymetry model of Smith and Sandwell's (SAS) approach (1994), GGM was found to have an advantage with short wavelength (≤ 12 km) components, while SAS better predicts longer wavelength (>25 km) components, despite its dependency on density contrast. To mitigate this limitation, a tuning density contrast of 10.25 g/cm³ between seawater and the seafloor was primarily estimated by the downward continuation method and then validated by the check points method with GGM. Similarly, SAS is limited by the "A" value in low-pass part of the Wiener filter, which defines the effective range of the wavelength components on bathymetry. As a final result, we present an enhanced GGM bathymetry model by integrating all available data.

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Introduction

Seafloor topography, or bathymetry, has traditionally been mapped by shipborne echo sounding measurements that are time-consuming and, because of the limited and uneven distribution of the measurements, often biased. However, the advent of the satellite radar altimeter has made it possible to estimate global bathymetry more economically and accurately. Since the late 1970 s, several satellite radar altimeters, such as Seasat (Born et al. 1979), Geosat (Cheney et al. 1986), ERS-1 and 2 (Gottschalk 1991; Francis et al. 1995), and TOPEX/Poseidon (Fu et al. 1994) have provided dense sea surface topography by measuring the distance between the satellite and the sea surface.

These measurements can be converted to geoid, free-air gravity anomalies (FAGA) and bathymetry. Since the variations of FAGA are theoretically correlated with undulations of crustal density variations of the local bedrock beneath the ocean floor, dense FAGA data that have been derived from satellite radar altimeter measurements can be used to predict bathymetry with the assumption that there is no horizontal density variation. Because of this limitation, shipborne bathymetry is required as a constraint.

Several methodologies for recovering bathymetry have been studied in recent decades. Dixon et al. (1983) compared the predicted bathymetry by one-dimensional filtering of the geoid heights obtained from tracks of the Seasat satellite altimeter (Born et al. 1979) with observed

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conventional bathymetry in the along-track directions. This study was followed by other one-dimensional along-track methods (Baudry et al. 1987; Craig and Sandwell 1988; Vogt and Jung 1991; Jung and Vogt 1992) and by two-dimensional bathymetry approaches (Baudry and Calmant 1991; Smith and Sandwell 1994; Calmant 1994; Sichoix and Bonneville 1996; Ramillien and Cazenave 1997; Hwang 1999; Calmant et al. 2002).

Calmant and Baudry (1996) described a comprehensive summary for bathymetry modeling of deep seafloor features, such as seamounts and ridges, on a global scale from satellite altimetry-derived FAGA. Bathymetric estimates derived from satellite altimetry and shipboard echo sounding data were reviewed by Sandwell and Smith (2001).

One of the useful techniques for estimating bathymetry is the gravity-geologic method (GGM). Originally, GGM was applied for mapping bedrock elevations beneath glacial drift deposits (Ibrahim and Hinze 1972; Adams and Hinze 1990; Nagarajan 1994). In terrestrial applications, GGM has been investigated for estimating the depth to bedrock from observed Bouguer gravity anomalies (BGA) and borehole measurements between the surface and the basement at the control points. One limitation in terrestrial application is that GGM predictions could be distorted because density varies between the glacial sediments and the bedrock. GGM has also been applied in marine settings, where the density contrast between seawater and the seafloor bedrock is very homogeneous. The method can be used for predicting the ocean bottom topography by using FAGA and depth measurements instead of BGA and borehole measurements.

In the past 40 years, several research studies have been performed for bathymetry predictions using GGM from satellite altimetry-derived FAGA. GGM was adapted for bathymetric estimations in the Barents Sea and around Greenland (Roman 1999). To implement GGM, a density contrast between seawater and the ocean bottom rock should be estimated. A density contrast of 9.0 g/cm³ was estimated for GGM and used in the Drake Passage, Antarctica (Kim et al. 2010). Determination of the effective density contrast estimated by the downward continuation method was recently implemented to predict bathymetry by GGM with an accuracy of less than 40 m (Hsiao et al. 2010).

In this study, combined shipborne and satellite altimetry measurements were used to estimate a 1×1 arc-minute bathymetry model, by using dense altimetry-derived FAGA where sparse shipborne depth measurements were available. The study area (36–39°N and 130–133°E) is located in the East Sea (Sea of Japan) between Korea and Japan (Fig. 1). As shown in Fig. 1, the Ulleung Basin is located south of Ulleung and Dok Islands.

Data used in this study included shipborne depths and gravity data from the National Geophysical Data Center



Fig. 1 Location map. Study area $(36-39^{\circ}N \text{ and } 130-133^{\circ}E)$ is delineated by a *red box*. The blue box is the boundary of the test area used for the density contrast determination. The 1×1 arc-minute bathymetry by ETOPO1 is superimposed as a background

(NGDC, http://www.ngdc.noaa.gov) GEODAS database (National Geophysical Data Center 2008), the National Oceanic and Atmospheric Administration (NOAA, http:// www.noaa.gov) and 1×1 arc-minute altimetry-derived FAGA (Sandwell and Smith 2009) generated by Geosat and ERS-1 satellite altimeters. Fig. 2a shows the distribution of the shipborne depths obtained from NGDC and the 1×1 arc-minute Earth topographical database 1 (ETO-PO1) model (Amante and Eakins 2009) as the background. In the left panel of Fig. 2a, white dots (18,317 points) have only shipborne depth measurements, and black dots (6,736 points) contain both depth and gravity measurements.

Figure 2b shows the altimetry-derived FAGA in the study area, which was obtained from Scripps Institution of Oceanography (SIO, http://www.sio.ucsd.edu), University of California at San Diego (Sandwell and Smith 2009). The original data were re-gridded to 1×1 arc-minutes by interpolation with the Generic Mapping Tools (GMT) software package (Wessel and Smith 1998), because there were the distortions in the high latitudes that were made when the map was generated from the Mercator projection.

To determine the density contrast, we first estimated a density contrast value by the downward continuation method and then evaluated to determine if the value was within the acceptable tuning density contrast range by GGM using check points (Hsiao et al. 2010). We also adopted Smith and Sandwell's (SAS) method (1994) to estimate bathymetry from FAGA in this study. The two bathymetry models were compared with other bathymetry models, including the NGDC shipborne data and the ETOPO1 model. Furthermore, we validated our bathymetry models

Fig. 2 a The 1×1 arc-minute bathymetry by ETOPO1 with superimposed ship track measurements by the National Geophysical Data Center (NGDC). Black dots denote points with depth and gravity, while white dots have depth only. The test area for the determination of density contrast is also shown (with control points as white circles and check points as black triangles). b Satellite altimetryderived free-air gravity anomalies (FAGA) (Sandwell and Smith 2009) in the study area. Attributes listed for this and subsequent maps include the amplitude range (AR = minimum and maximum values), amplitude mean (AM), amplitude standard deviation (ASD), and amplitude unit (AU)

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AR=(-41.2, 122.2), AM=4.5, ASD=20.6, AU=mGal

with independently measured shipborne data by the Korea Institute of Geoscience and Mineral Resources (KIGAM, http://www.kigam.re.kr) in April 2000. In this paper, we also present an enhanced bathymetry model by integrating all shipborne and satellite measurements in the study area as the final result of the study.

Bathymetry prediction by the gravity-geologic method

Concept of the gravity-geologic method

GGM was originally developed to predict the depth-tobasement using a density contrast between glacial sediments and bedrock from known gravity anomalies and borehole measurements (Ibrahim and Hinze 1972; Adams and Hinze 1990; Nagarajan 1994).

In general, the observed BGA includes two gravity effects that are composed of the residual gravity generated from local bedrock variations and the regional gravity generated from deeper mass variations. In GGM applications, we need the residual gravity effect, which yields short wavelength signals, to predict the final depth-tobedrock. The regional gravity is computed by subtracting the residual gravity from the observed BGA if the depth measurements are known at the control points. However, at sites where the depth is not measured, we can estimate the residual gravity for predicting the bedrock undulations by subtracting the regional gravity, which is generated by gridding the regional gravity yielded at sites of measured depth, j, from the observed BGA. The control points with measured depth are denoted with subscript, j, while the control points with unmeasured depth are indicated with subscript i in this study.

Local bedrock variations under sediments yield a shorter wavelength gravity field, while deeper mass variations generate a longer wavelength gravity field. These two gravity fields comprise the observed Bouguer gravity (g_{OBS}) by combining the residual gravity (g_{RES}) and the regional gravity (g_{REG}) as in the following equation:

$$g_{OBS}(i) = g_{RES}(i) + g_{REG}(i).$$
(1)

We assume that the measured depths at control points j_k in Fig. 3 were used to estimate the residual bedrock gravity field (g_{RES}) that generates the shorter wavelength effect from a simple Bouguer slab formula:

$$g_{RES}(j) = 2\pi G(\Delta \rho)(E(j) - D), \qquad (2)$$

where G is the gravitational constant, $6.672 \times 10^{-8} \text{ cm}^3/\text{g} \cdot \text{sec}^2$; $\Delta \rho$ is the density contrast (g/cm³); E(j) is the bedrock elevation at the *j*th control point; and, D is the deepest depth of the control points as a reference datum. Both E(j) and D are measured in meters.

The residual gravity represents the effect of the bedrock surface at the control points and is removed from the observed Bouguer gravity to estimate the regional gravity field (g_{REG}) that represents the longer wavelength effect:

$$g_{REG}(j) = g_{OBS}(j) - g_{RES}(j).$$
(3)

The regional gravity, $g_{REG}(i)$, at sites of unmeasured depth, *i*, can be estimated by interpolation as a grid by using the generated regional gravity, $g_{REG}(j)$, at sites of measured depth, *j*. The estimated regional gravity, $g_{REG}(i)$, at site *i* is removed from the observed gravity, $g_{OBS}(i)$, to estimate the residual gravity, $g_{RES}(i)$:

$$g_{RES}(i) = g_{OBS}(i) - g_{REG}(i).$$
(4)



Fig. 3 Geometry of the gravity-geologic method and the downward continuation method (modified from Hsiao et al. 2010)

By rearranging Eq. 2, the elevation of the bedrock above the datum where bedrock depths are unmeasured at site i can be estimated with the following formula:

$$E(i) = \frac{g_{RES}(i)}{2\pi G \Delta \rho} + D.$$
(5)

The estimation of bedrock topography using GGM is influenced by the distribution and number of the depths and gravity at the control points, the observed gravity, and the density contrast between the bedrock and the overlying sediments. Where the shipborne depth measurements are sparse over oceanic areas, GGM can be used to predict bathymetry from satellite altimetry-derived dense FAGA. GGM applications in ocean areas use shipborne depths and FAGA, instead of bedrock elevation and Bouguer gravity at control points, in order to predict the ocean bottom topography under seawater (Roman 1999). FAGA data are approximately at the geoid in the ocean and, thus, do not need correction for height variation and Bouguer mass effects.

The density contrast between seawater and the seafloor varies as a result of changes in mass under the ocean bottom. In GGM applications, a geologically reasonable density contrast between seawater (1.03 g/cm^3) and the seafloor bedrock (2.70 g/cm^3) is assumed to be 1.67 g/cm³. The density of the seafloor bedrock in the crust is assumed to be the mean value between 2.67 g/cm³ and 2.73 g/cm³ (Jin 1995; Kim et al. 2010). In reality, the density contrast is much larger than the theoretical value (1.67 g/cm^3) , because satellite altimetry-derived FAGA compiled on the sea surface in GGM are caused by the sources at the ocean bottom (Roman 1999; Kim et al. 2010). For more effective bathymetric prediction, a tuning value of density contrast that was determined from the control points of measured shipborne depths was used (Kim et al. 2010; Hsiao et al. 2010).

Determination of density contrast

Check points with the gravity-geologic method

Because the shipborne depth measurements shown in Fig. 2a have large gaps, we selected the southeast corner (36.0–37.46°N and 132.14–133.0°E) as a test area for estimating the density contrast. This test area is delineated by the blue box in Figs. 1 and 2a, where the distribution of shipborne measurements is denser and more even. The right-hand panel in Fig. 2a shows control and check points within the test area with the ETOPO1 bathymetry model as the background. Of all 1,493 shipborne depth measurements within the test area, two thirds (992) of the points were used as control points (white circles). The remaining third (501) were used as check points (black triangles). Every third point of shipborne depths along shiptracks was

selected as a check point to evenly distribute the control points within the test area. The control points in the test area were used to effectively determine the density contrast of the study area.

Figure 4 is a trade-off diagram, which shows a rootmean-square (RMS) difference, its rate of change and a correlation coefficient between the control and check points. This trade-off diagram was developed by the check points method for picking an effective density contrast for GGM bathymetry predictions. If the selected density contrast is too small, the GGM-predicted bathymetry would be very large. In contrast, if the selected density contrast is too big, the GGM-predicted bathymetry would be too close to "level n" (Fig. 3). A range of density contrasts that stabilize the RMS difference by using the check points method was selected in the trade-off diagram. From the three curves in the trade-off diagram, we selected the range of acceptable density contrasts as 9.0 g/cm³ and higher, which is the shaded area in Fig. 4.

The determination of density contrast between seawater and the ocean bottom topographic mass is an important factor for GGM bathymetry. We know that the density contrast ($\Delta \rho$) in Eq. 5 controls the amplitude of bathymetric estimates. Although it is not geologically reasonable, the check points were used to select an optimal density contrast in the GGM application (Nagarajan 1994; Kim et al. 2010).

Downward continuation method

In this study, the downward continuation method was used to select an effective density contrast within the acceptable range, i.e. 9.0 g/cm^3 and higher, determined by the check points method described in the previous section. Upward continuation is regarded as a smooth operation (Tikhonov 289

and Arsenin 1977), which is a well-posed problem whose solution will be unique; whereas, downward continuation is characterized as an unstable solution, which is an ill-posed problem whose solution will be non-unique and unstable. Upward continuation in the frequency domain from the gravity field at $z = h_1$ plane to $z = h_2$ plane can be explained as (Parker 1977):

$$G_{h_2}(u,v) = \exp(-2\pi kd)G_{h_1}(u,v),$$
(6)

where $G_{h_1}(u, v)$ and $G_{h_2}(u, v)$ are the two-dimensional Fourier transforms of the gravity field at h_1 and h_2 , respectively; u and v are the frequencies for x and ydirections, respectively; $k = \sqrt{u^2 + v^2}$; and $d = h_2 - h_1$. Downward continuation of the gravity field in the Fourier transform from $z = h_2$ plane to $z = h_1$ plane, by contrast, is given by inverting Eq. 6 as:

$$G_{h_1}(u, v) = \exp(2\pi kd)G_{h_2}(u, v).$$
(7)

The sea surface gravity field using the downward continuation method was calculated for several downward continued levels, as shown in Fig. 3, by the Gaussian filter in Eq. 7:

$$\bar{G}(u,v) = F(u,v)G(u,v)\exp(2\pi kd), \tag{8}$$

where \overline{G} represents the downward continued gravity field, G represents the original gravity field, d denotes the distance of the downward continuation, and F is a Gaussian filter.

The effective density contrast for considering mass variations in the deepest ocean bottom was determined by the downward continuation method until reaching the seafloor topography. The density contrast from the estimated gravity ratio between the sea surface level and each downward continued level was computed, as shown in Fig. 3. The density contrast ($\Delta \rho_{\text{DWC}}$) at point *i* of every

Fig. 4 Trade-off diagram for selecting the tuning density contrast range of 9.0 g/cm³ and higher by the gravity-geologic method and a final density contrast of 10.25 g/cm³ by the downward continuation method



downward continued level can be estimated as a multiplying gravity ratio, $\Delta g_{\text{DWC}}(i)/\Delta g_{\text{SEA}}(i)$, at point *i* of each downward continued level and sea surface level by $\Delta \rho_{\text{SEA}}$, which is the density contrast at the sea surface level as a constant of 1.03 g/cm³ (Strykowski et al. 2005). The density contrast in each downward continued level was computed as the mean value of $\Delta \rho_{\text{DWC}}$.

Figure 5 shows the predicted density contrasts for every level at intervals of 200 m below the sea surface, along with the gravity ratios between the sea surface and the downward continued level. The density contrast of each level was calculated until the levels reached 2,037 m, which is the greatest depth in the test area. As shown in Fig. 5, both the gravity ratio and the predicted density contrast gradually increased with the ocean depth.

From the two curves, 10.25 g/cm³, which is the predicted density contrast at the deepest depth, was selected as a tuning density contrast that would be applied to the GGM bathymetry predictions, because the tuning density contrast is within the acceptable range determined by the check points method with GGM.

Bathymetry prediction by Smith and Sandwell's approach

Review of Smith and Sandwell's approach

The Parker formula (Parker 1972) can be described as a forward model used to convert seafloor topography (*h*) to gravity anomaly (Δg) in the two-dimensional Fourier domain:



Fig. 5 Density contrast at every level and gravity ratio between the sea surface and the downward continued level in the test area

$$\Im[\Delta g] = 2\pi G(\Delta \rho) \exp(-2\pi kd) \sum_{n=1}^{\infty} \frac{(2\pi k)^{n-1}}{n!} \Im[h^n], \qquad (9)$$

where *G* is the gravitational constant; $\Delta \rho$ is the density contrast between seawater and the seafloor topography; *d* is the mean depth beneath the ocean surface; *k* is the radial frequency $\left(=\sqrt{k_x^2 + k_y^2}\right)$; *k* is $(k_x, k_y) = (1/\lambda_x, 1/\lambda_y)$, where (k_x, k_y) and (λ_x, λ_y) are frequencies and wavelengths in the *x* and *y* directions, respectively; and, \Im [] is the Fourier transform operator.

If the magnitude of the seafloor topography undulation is much less than the mean ocean depth (d), the first term dominates because the series of non-linear terms of Eq. 9 converge fast. The first term in Eq. 9 is only dependent on the mean depth below the ocean surface and is a linear relationship between the gravity anomaly and the seafloor topography in the Fourier domain as in the following equation:

$$G(\mathbf{k}) = Z(k)H(\mathbf{k}) = 2\pi G(\Delta \rho) \exp(-2\pi kd)H(\mathbf{k}), \qquad (10)$$

where G(k) and H(k) are the Fourier transforms of the gravity anomalies (Δg) and the seafloor topography (h), respectively; Z(k) is the admittance function as a transfer function; and, k is the frequency in two-dimensions. Z(k) is a transfer function that is isotropic and spatially invariant.

In the Z(k) terms of Eq. 10, the $2\pi G\Delta\rho$ term is a Bouguer constant as an infinite slab of material; and, the exp($-2\pi kd$) term represents upward continuation (from seafloor to sea surface) through the mean ocean depth as exponential decay with increasing frequency.

Next, we estimated the seafloor topography from satellite altimetry-derived gravity anomalies using the inverse model of Eq. 10. The inversion model is shown as the following equation in the Fourier domain:

$$H(\boldsymbol{k}) = Z^{-1}(k)G(\boldsymbol{k}) = \frac{1}{2\pi G\Delta\rho} \cdot \exp(2\pi kd)G(\boldsymbol{k}).$$
(11)

Smith and Sandwell (1994) (SAS) suggested that, in a limited band of intermediate wavelengths between about 15 and 160 km, the satellite altimetry-derived FAGA data have a correlation between gravity and bathymetry, because they can be implemented by attenuation of gravity by upward continuation in shorter wavelengths and by isostatic compensation of the seafloor topography at longer wavelengths. For a stable condition of downward continuation, which is represented as the term $\exp(2\pi kd)$ in Eq. 11, to estimate bathymetry from gravity anomalies, the transfer function $Z^{-1}(k)$ is suppressed by windowing the limited range of wavelengths of topography. Thus, a bandpass filter was used. The purpose of this filter is to stabilize the downward continuation in predictions of the seafloor topography (Smith and Sandwell 1994; Calmant and Baudry 1996; Watts et al. 2006).

In the 1994 SAS approach, the altimetry-derived gravity anomalies in the frequency domain, $G_0(\mathbf{k})$, are first bandpass filtered and downward continued to a mean depth (d) by:

$$G(\mathbf{k}) = G_0(\mathbf{k})W(k)\exp(2\pi kd), \qquad (12)$$

where $G(\mathbf{k})$ is the band-pass filtered and downward continued gravity field, and W(k) is the Wiener filter. The gridded shipborne depth measurements in the frequency domain, $B_0(\mathbf{k})$, are also band-pass filtered by the Wiener filter:

$$H(\boldsymbol{k}) = B_0(\boldsymbol{k})W(\boldsymbol{k}),\tag{13}$$

where $H(\mathbf{k})$ is the seafloor topography field.

In Eqs. 12 and 13, the Wiener filter is composed of two filters, $W_1(k)$ as the high-pass filter and $W_2(k)$ as the low-pass filter:

$$W(k) = W_1(k) \cdot W_2(k),$$
 (14)

where $W_1(k) = 1 - \exp[-2(\pi ks)^2]$, $W_2(k) = [1 + Ak^4 \exp(4\pi kd)]^{-1}$, and *s* is the Wiener filter parameter.

The "A" value chosen in $W_2(k)$ affects the cut-off wavelength scale at the resolution of the predicted bathymetry. Thus, the choice of the "A" value must be considered rigorously and should represent an appropriate compromise between noise reduction and signal detail. Unlike SAS (1994), who used 9,500 km⁴ for the "A" in the low-pass filtering (i.e. $W_2(k)$), we selected 6,233 km⁴ in our study with a new gravity model (personal communication with David Sandwell, 2009). The gridded shipborne depth measurements in the frequency domain, $B_0(\mathbf{k})$, are then low-pass filtered by:

$$D_{LP}(\mathbf{k}) = B_0(\mathbf{k})[1 - W_1(k)], \tag{15}$$

where $D_{LP}(k)$ is the low-pass filtered gridded shipborne depth in the frequency domain.

Finally, the SAS bathymetry is estimated by:

$$SAS = d_{LP}(x) + S \cdot g(x), \tag{16}$$

where $d_{LP}(x)$ and g(x) are the spatial domains of $D_{LP}(\mathbf{k})$ in Eq. 15 and $G(\mathbf{k})$ in Eq. 12, respectively; S is a scaling factor determined by linear regression between $G(\mathbf{k})$ in Eq. 12 and $H(\mathbf{k})$ in Eq. 13; and, $d_{LP}(x)$ and $S \cdot g(x)$ represent the longer and band-limited wavelength bathymetry, respectively.

Results from Smith and Sandwell's approach

The SAS bathymetry model is shown in Fig. 6b. In SAS, the linear relationship between gravity anomalies and the seafloor topography was fundamentally used in the Fourier domain. Band-pass filtering was performed in the downward continued satellite altimetry gravity anomalies and the gridded shipborne depths using Eqs. 12 and 13,

respectively, with a mean water depth of 1,868 m in the study area. g(x) in the spatial domain obtained from Eq. 12 was converted into the seafloor topography by multiplication with the scaling factor, *S*, which was empirically computed from linear regressions of band-pass filtered and downward continued gravity anomalies, g(x), and band-pass filtered gridded shipborne depths, h(x), in the study area.

The scaling factor of 2.20 m/mGal was used to convert band-pass filtered and downward continued altimetryderived gravity anomalies into bathymetry. For longer wavelengths with regional effect, low-pass filtering of the gridded shipborne depth measurements was implemented in Eq. 15. Finally, the SAS bathymetry model was generated from two bathymetries at different wavelengths.

Comparison of bathymetry models

The 1×1 arc-minute bathymetry models by GGM, SAS, NGDC, and ETOPO1 are shown in Fig. 6. The NGDC gridded model was produced by the harmonic surface technique in GMT software (Wessel and Smith 1998). Statistical comparisons are summarized in Table 1. The GGM model showed correlation coefficients of 0.998, 0.999, and 0.957 with SAS, NGDC, and ETOPO1 models, respectively. SAS also had a maximum correlation coefficient of 0.998 with NGDC.

The reason why both GGM and SAS were most strongly correlated with the NGDC gridded model is that NGDC shipborne depths were utilized to estimate both bathymetry models. According to the comparison in Table 1, the RMS of the GGM model is smaller than those of the SAS, NGDC and ETOPO1 models.

Maps of point-by-point differences between the models are represented in Fig. 7, and Table 2 summarizes their statistics. In Fig. 7b, the difference between the GGM and NGDC gridded models (GGM–NGDC) represents good agreement with the smallest RMS (34.1 m) in Table 2. However, the differences in shallow water around Ulleung Island in Fig. 7c, e, f reached approximately -2,000 m. These large differences were caused by the steep bathymetry gradient around the island.

The power spectral density of the bathymetry models was analyzed in the radial direction to examine energy variations in different wavelengths. In Fig. 8, the amplitude of the power spectral density along the y-axis is represented as a unit of dB that is computed from $10 \cdot \log_{10}(P)$, where P is a power. The wavelength along the x-axis is considered to range up to 300 km due to the size of the study area. As illustrated in Fig. 8, the SAS model showed a high power spectral density for wavelengths longer than

Fig. 6 Bathymetry models by a the gravity-geologic method (GGM), b the Smith and Sandwell (SAS) approach (1994), c the National Geophysical Data Center (NGDC) and d Earth topographical database 1 (ETOPO1). Bathymetry estimates along Profiles #1 and #2 are contrasted. *Red dots* denote shipborne measurements with anomalous depths



(b) SAS



AR=(-3998.1, -49.8), AM=-1853.4, ASD=693.7, AU=meter





 Table 1
 Statistic comparisons between bathymetry models by the gravity-geologic method (GGM), Smith and Sandwell's (SAS) approach (1994), National Geophysical Data Center (NGDC), and Earth topographical database 1 (ETOPO1) in the study area (unit: meter)

	Min.	Max.	Mean	SD	RMS	Correlatio	Correlation coefficient (CC)		
						GGM	SAS	NGDC	ETOPO1
GGM	-3998.1	-49.8	-1853.4	693.7	1979.0	1	0.998	0.999	0.957
SAS	-3238.8	0.0	-1853.4	720.7	1988.6	0.998	1	0.998	0.959
NGDC	-3238.6	-52.8	-1857.7	693.4	1982.9	0.999	0.998	1	0.950
ETOPO1	-3427.0	0.0	-1849.4	707.0	1979.9	0.957	0.959	0.950	1

25 km and presented low energy in short wavelengths less than roughly 12 km.

Both the GGM and SAS models showed similar energy levels in the medium wavelengths between 12 km and 25 km; however, the GGM model had more power in very short wavelengths of less than approximately 12 km. These results may indicate that the GGM has better bathymetry in the short wavelength components; whereas, the SAS can effectively yield bathymetry in the medium to longer wavelength components.

Fig. 7 The differences between four bathymetry models: GGM, SAS, NGDC and ETOPO1. a GGM – SAS, b GGM – NGDC, c GGM – ETOPO1, d SAS – NGDC, e SAS – ETOPO1, and f NGDC – ETOPO1



To validate the result of the power spectral density analysis, we directly compared the GGM and SAS models with NGDC shipborne measurements possibly with short wavelength components. Before the comparison, the two models were interpolated into the NGDC shipborne locations, which are shown as black dots in Fig. 2a. Figure 9a,

	Min.	Max.	Mean	SD	RMS
GGM – SAS	-1665.3	304.6	0.0	52.8	52.8
GGM – NGDC	-1596.2	289.1	4.3	33.8	34.1
GGM – ETOPO1	-1877.7	967.3	-4.0	206.6	206.6
SAS – NGDC	-119.5	372.4	4.3	53.6	53.8
SAS – ETOPO1	-1801.1	934.6	-4.1	205.8	205.9
NGDC – ETOPO1	-2128.4	946.0	-8.4	221.3	221.5

 Table 2 Statistical comparison of the differences between four

 bathymetry models: GGM, SAS, NGDC, and ETOPO1 (unit: meter)



Fig. 8 The comparison of power spectral densities between GGM (*red circles*) and SAS (*blue circles*) bathymetry models. The unit of power spectral density is dB, which is computed from 10-log(P), where P is a power

b show the absolute values of the differences between the NGDC ship depths and the GGM bathymetry and between the NGDC ship depths and the SAS bathymetry on the NGDC shipborne tracks.

As shown in Fig. 9a, the minimum, maximum, mean, and standard deviation of the absolute differences are 0.0, 1222.7, 16.1, and 37.8 m, respectively. Ninety-three percent of the absolute differences were less than 37.8 m, which is the standard deviation. Only 1.5% of the absolute differences were more than 100 m in the NGDC ship tracks in the study area. In Fig. 9b, the statistics of the absolute differences between the SAS bathymetry and the NGDC bathymetry on the NGDC ship tracks were 0.0, 656.7, 46.9, and 44.4 m for the minimum, maximum, mean, and standard deviation, respectively. The absolute difference below 44.4 m (the standard deviation) was 60.5% and 10.9% for more than 100 m.

From the histograms of two absolute differences from the NGDC ship tracks and the two competing models presented in Fig. 9c, d, we know that 95% of the absolute differences between NGDC and GGM were below 50.0 m; whereas, 67.4% of the absolute differences between NGDC and SAS were less than 50.0 m. Based on the small differences between the GGM model and the NGDC shipborne depths along the NGDC ship tracks, we concluded that GGM is more suitable than the SAS method for short wavelength bathymetry prediction.

GGM bathymetry estimated using a tuning density contrast (10.25 g/cm³) and determined by the downward continuation method is enhanced with an accuracy of 28.4 m in the standard deviation (693.7 m in Table 1), when comparing the standard deviation (722.1 m) of the GGM bathymetry predicted by the geologically reasonable density contrast (1.67 g/cm³) between seawater and the seafloor topography.

In addition to the NGDC shipborne depths, shipborne data measured in April 2000 by the Korea Institute of Geoscience and Mineral Resources (KIGAM), which are not included in the NGDC, NOAA databases, were compared with the GGM, SAS, and NGDC models along two profiles: Profile #1 (east-west direction) and Profile #2 (north-south direction), shown in white lines in Fig. 6a, b, c. The results are shown in Fig. 10.

The KIGAM shipborne depths in Profile #1 were strongly correlated with GGM, SAS, and NGDC with correlation coefficients of 0.982, 0.976, and 0.979, respectively. For Profile #2, the correlation coefficients were also very strong with 0.997, 0.998, and 0.998 for GGM, SAS, and NGDC, respectively. Table 3 summarizes the statistics of KIGAM, GGM, SAS, and NGDC depths along Profiles #1 and #2 of the KIGAM ship tracks.

Table 4 presents the statistics of the differences between KIGAM and GGM, SAS, and NGDC bathymetries along Profiles #1 and #2. In Profile #1, the RMS of the differences between the KIGAM ship bathymetry and each of the GGM, SAS, and NGDC bathymetries were 43.5, 79.7 and 46.8 m, respectively; and, in Profile #2, the difference between KIGAM and GGM had the smallest RMS value (28.7 m).

However, because of the overestimated SAS bathymetry in shallow water areas in Fig. 10a, b, the profiles between KIGAM and SAS had a discrepancy of about 100 m between 130.45 and 130.65°E in Profile #1 and between 37.65 and 37.72°N in Profile #2. These differences were caused by the scaling factor (*S*) in the SAS bathymetry model; therefore, we can decrease this difference by adjusting *S*. We know that short wavelength bathymetry is well recovered by the GGM model, because the GGM bathymetry indicated small RMS differences compared with the two profiles (Profiles #1 and #2) collected along



Fig. 9 The absolute differences of bathymetry on the NGDC shipborne tracks **a** between NGDC and GGM and **b** between NGDC and SAS. The histograms of the absolute differences on the NGDC shipborne tracks **c** NGDC and GGM and **d** NGDC and SAS

the KIGAM ship tracks with very short wavelength bathymetry.

After examination of the maps in Fig. 6, we found that the three bathymetry models (GGM, SAS, and NGDC) had ship tracks with outlier data, when compared with the ETOPO1 model. The outliers were considered as anomalous depths and are shown as red dots in Fig. 6a–d. We eliminated ship tracks with outliers, as shown in red dots in the left panel of Fig. 2a. The eliminated shipborne depths in Fig. 2a were not used for predicting an enhanced bathymetry model.

As a final result, we present an enhanced bathymetry model by using all shipborne measurements, including KIGAM data, in GGM, with a density contrast of 10.25 g/ cm^3 estimated by GGM and the downward continuation method in the test area (Fig. 11a). The ETOPO1 model was also applied to the islands, where the topography is above sea level. The enhanced bathymetry model had a minimum, maximum, mean, and standard deviation of -3182.3, 756.0, -1858.9, and 699.7 m, respectively.

The white color in the new model shown in Fig. 11a denotes depth at or above sea level. Figure 11b shows the difference between the enhanced and original GGM models. Overall, the enhanced model was similar to the original GGM model in Fig. 6a, except for the shallow water around Ulleung and Dok Islands, where the new model



Fig. 10 Bathymetry estimate comparisons for Profiles #1 (a) and #2 (b), respectively, as shown in Fig. 6. CC is correlation coefficient

Table 3 Statistic comparisons between the Korea Institute of Geo-science and Mineral Resources (KIGAM), GGM, SAS, and NGDCbathymetry along Profiles #1 and #2 (unit: meter)

		Min.	Max.	Mean	SD	RMS
KIGAM	Profile #1	-2032.5	-1225.2	-1487.5	185.9	1499.0
	Profile #2	-2080.8	-1235.4	-1611.5	323.2	1643.4
GGM	Profile #1	-1908.7	-1222.9	-1461.8	185.2	1473.4
	Profile #2	-2078.0	-1221.1	-1610.0	309.7	1639.3
SAS	Profile #1	-1949.9	-1133.7	-1442.1	233.0	1460.7
	Profile #2	-2087.9	-1145.9	-1563.1	341.4	1599.7
NGDC	Profile #1	-1931.1	-1216.5	-1464.1	197.0	1477.2
	Profile #2	-2070.3	-1214.8	-1597.5	306.6	1626.5

 Table 4
 Statistics of the differences between KIGAM, GGM, SAS, and NGDC bathymetry in Profiles #1 and #2 (unit: meter)

		Min.	Max.	Mean	SD	RMS
KIGAM – GGM	Profile #1	-123.8	74.6	-25.6	35.2	43.5
	Profile #2	-62.7	57.4	-1.5	28.8	28.7
KIGAM – SAS	Profile #1	-137.4	103.5	-45.4	65.8	79.7
	Profile #2	-94.4	7.1	-48.5	27.2	55.5
KIGAM – NGDC	Profile #1	-101.4	94.1	-23.4	40.7	46.8
	Profile #2	-73.5	38.5	-14.0	26.7	30.1

provided a better estimation by ETOPO1 for the area above sea level. The correlation coefficient between the enhanced and original GGM models was 0.966.

Results and conclusions

In this study, the gravity-geologic method (GGM) was used to enhance a bathymetry model from NGDC shipborne depths and satellite altimetry-derived FAGA (Sandwell and Smith 2009) in the area of 36–39° N and 130–133° E in the East Sea. The denser shipborne measurements of the test area in the southeast part of the study area were used to calculate the density contrasts between the seawater and the ocean bottom topography. A tuning density contrast of 10.25 g/cm³, which was estimated by the downward continuation method, was found to be in the acceptable range (9.0 g/cm³ and higher), as determined with control and check points.

According to the flexural isostatic compensation theory, the relationship between bathymetry and downward continued gravity anomalies is correlated in the 15–160 km wavelength band (Smith and Sandwell 1994). We implemented the Smith and Sandwell's (SAS) approach (1994) by using the linear relationship between the seafloor topography and the downward continued gravity anomalies in the study area. The GGM bathymetry had correlation coefficients of 0.998, 0.999, and 0.957 with the SAS, NGDC, and ETOPO1 bathymetry models, respectively. The difference between the GGM and NGDC gridded models from the results of comparisons of the four bathymetry models (GGM, SAS, NGDC gridded, and ETOPO1) presented smaller RMS than the differences with the others.

In addition, the standard deviation of the absolute differences between the GGM bathymetry model and the NGDC shipborne depths on NGDC ship tracks was below 37.8 m. From the power spectral density analysis, we showed that GGM has an advantage with short wavelength (\leq 12 km) components, while the SAS provides a better prediction of longer wavelength (\geq 25 km) components. In the medium wavelength range, the two methods perform similarly.

An enhanced GGM bathymetry model was presented as a final result by combining all available data.

We conclude that GGM better predicts bathymetry over short wavelengths (≤ 12 km) than SAS. Moreover, the bathymetry calculated by GGM is easily upgradable with



Fig. 11 a Enhanced bathymetry model and b differenced (Fig. 11a – Fig. 6a) GGM bathymetry models. Note that the two *color scales* are different

future satellite and shipborne measurements. Because the accuracy of GGM is dependant on the density contrast, we present a method for estimating a suitable tuning density contrast by combining the GGM and downward continuation methods. The use of the Bouguer slab approximation mitigated by establishing tuning densities stabilized the GGM predictions (Kim et al. 2010), especially for modeling the gravity effects of variable density and rugged bathymetric relief at distances up to several kilometers.

In fact, the estimated tuning density of 10.25 g/cm³ was found to have a more analytical than physical significance, because it essentially mitigates limitations in using a single-density Bouguer slab for analyzing multi-density terrain with rugged relief components that require gravity terrain corrections (Nagarajan 1994; Kim et al. 2010). The GGM bathymetry predicted by the tuning density contrast (10.25 g/cm^3) presents an improvement of 28.4 m over the GGM bathymetry estimated by the reference density contrast (1.67 g/cm³) between seawater and the ocean bottom topography, although the tuning density is larger than a geologically reasonable value. The GGM model using the geologically reasonable density contrast is not realistic. It implies that the influence of density contrast for GGM is unavoidable. Because the SAS approach includes a bandpass filter, it may result in underestimating the short wavelength components.

Geology and bathymetry of the East Sea recently became an important issue, not only because understanding the origin of the basins in the East Sea is important for evaluation of natural resources (i.e. gas hydrate and hydrocarbon), but also in terms of exclusive economic zones (EEZs) and territorial boundaries between the neighboring countries. As a result, the GGM method is a useful and convenient approach for updating bathymetry when new satellite altimetry mission gravity data and shipborne depths become available.

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